

# Uncertainty relations with the variance and the quantum Fisher information

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## 1 Motivation

- How can we improve uncertainty relations?

## 2 Background

- Quantum Fisher information
- Recent findings on the quantum Fisher information
- Uncertainty relations

## 3 Uncertainty relations with the variance and the QFI

- Uncertainty relations based on a convex roof of the bound
- Cramér-Rao bound based on a convex roof
- Uncertainty relations based on a concave roof of the bound
- Uncertainty relations with several variances and the QFI
- Simple observation to prove further relations
- Metrological usefulness and entanglement conditions

# How can we improve uncertainty relations?

- There are many approaches to improve uncertainty relations.
- We show a method that replaces the variance with the quantum Fisher information in some well known uncertainty relations.
- We use convex/concave roofs over the decompositions of the density matrix.

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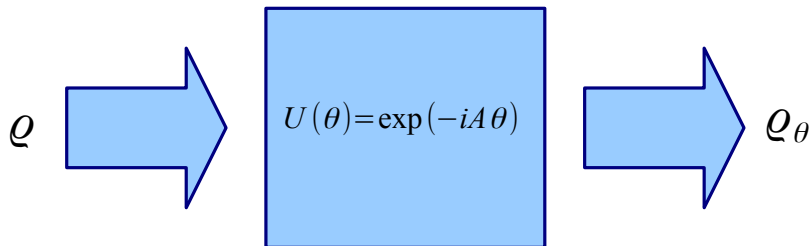
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# Quantum metrology

- Fundamental task in metrology



- We have to estimate  $\theta$  in the dynamics

$$U = \exp(-iA\theta).$$

# The quantum Fisher information

- Cramér-Rao bound on the precision of parameter estimation

$$(\Delta\theta)^2 \geq \frac{1}{mF_Q[\varrho, \mathbf{A}]},$$

where  $F_Q[\varrho, \mathbf{A}]$  is the **quantum Fisher information**, and  $m$  is the number of independent repetitions.

- The quantum Fisher information is

$$F_Q[\varrho, \mathbf{A}] = 2 \sum_{k,l} \frac{(\lambda_k - \lambda_l)^2}{\lambda_k + \lambda_l} |\langle k | \mathbf{A} | l \rangle|^2,$$

where  $\varrho = \sum_k \lambda_k |k\rangle\langle k|$ .

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# Properties of the Fisher information

Many bounds on the quantum Fisher information can be derived from these simple properties:

- For pure states, it equals four times the variance,  
$$F[|\psi\rangle\langle\psi|, A] = 4(\Delta A)^2_\psi.$$
- For mixed states, it is convex.



# The quantum Fisher information vs. entanglement

- For separable states

$$F_Q[\varrho, J_l] \leq N, \quad l = x, y, z.$$

[Pezze, Smerzi, Phys. Rev. Lett. 102, 100401 (2009); Hyllus, Gühne, Smerzi, Phys. Rev. A 82, 012337 (2010)]

- For states with at most  $k$ -particle entanglement ( $k$  is divisor of  $N$ )

$$F_Q[\varrho, J_l] \leq kN.$$

[P. Hyllus *et al.*, Phys. Rev. A 85, 022321 (2012); GT, Phys. Rev. A 85, 022322 (2012)].

- Macroscopic superpositions (e.g, GHZ states, Dicke states)

$$F_Q[\varrho, J_l] \propto N^2,$$

[F. Fröwis, W. Dür, New J. Phys. 14 093039 (2012).]

# Most important characteristics used for estimation

The quantum Fisher information is the convex roof of the variance

$$F_Q[\varrho, A] = 4 \min_{\{\rho_k, |\psi_k\rangle\}} \sum_k \rho_k (\Delta A)_{\psi_k}^2,$$

where

$$\varrho = \sum_k \rho_k |\psi_k\rangle \langle \psi_k|.$$

[GT, D. Petz, Phys. Rev. A 87, 032324 (2013); S. Yu, arXiv1302.5311 (2013);  
GT, I. Apellaniz, J. Phys. A: Math. Theor. 47, 424006 (2014)]

- Extended convexity for non-unitary dynamics.  
[S. Alipour, A. T. Rezakhani, Phys. Rev. A 91, 042104 (2015).]
- Convex roof over purifications.  
[R. Demkowicz-Dobrzański, J. Kołodyński, M. Guţă, Nature Comm. 2012.]

# Formula based on concave roofs

The variance is the concave roof of itself

$$(\Delta A)^2_{\varrho} = \max_{\{p_k, |\psi_k\rangle\}} \sum_k p_k (\Delta A)^2_{\psi_k},$$

where

$$\varrho = \sum_k p_k |\psi_k\rangle \langle \psi_k|.$$

GT, D. Petz, Phys. Rev. A 87, 032324 (2013);

GT, I. Apellaniz, J. Phys. A: Math. Theor. 47, 424006 (2014).

## A single relation for the QFI and the variance

The previous statements can be concisely reformulated as follows. For any decomposition  $\{p_k, |\psi_k\rangle\}$  of the density matrix  $\varrho$  we have

$$\frac{1}{4}F_Q[\varrho, A] \leq \sum_k p_k (\Delta A)_{\psi_k}^2 \leq (\Delta A)_{\varrho}^2,$$

where the upper and the lower bounds are both **tight**.

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# Robertson-Schrödinger inequality

The Robertson-Schrödinger inequality is defined as

$$(\Delta A)_\rho^2 (\Delta B)_\rho^2 \geq \frac{1}{4} |L_\rho|^2,$$

where the lower bound is given by

$$L_\rho = \sqrt{|\langle \{A, B\} \rangle_\rho - 2\langle A \rangle_\rho \langle B \rangle_\rho|^2 + |\langle C \rangle_\rho|^2},$$

$\{A, B\} = AB + BA$  is the anticommutator, and we used the definition

$$C = i[A, B].$$

Important:  $L_\rho$  is neither convex nor concave in  $\rho$ .

# Heisenberg uncertainty

The Heisenberg inequality is defined as

$$(\Delta A)_{\rho}^2 (\Delta B)_{\rho}^2 \geq \frac{1}{4} |\langle C \rangle_{\rho}|^2,$$

where we used the definition

$$C = i[A, B].$$

## The two inequalities together

We have two inequalities

$$(\Delta A)^2_{\rho}(\Delta B)^2_{\rho} \geq \frac{1}{4}|L_{\rho}|^2 \geq \frac{1}{4}|\langle C \rangle_{\rho}|^2.$$

The Heisenberg uncertainty can be saturated only if

$$|L_{\rho}| = |\langle C \rangle_{\rho}|.$$



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# New inequality from decompositions

- Consider a decomposition to mixed states

$$\varrho = \sum_k p_k \varrho_k.$$

- For such a decomposition, for all  $\varrho_k$  the Robertson-Schrödinger inequality holds

$$(\Delta A)_{\varrho_k}^2 (\Delta B)_{\varrho_k}^2 \geq \frac{1}{4} |L_{\varrho_k}|^2.$$

- Let us consider the inequality

$$\left( \sum_k p_k a_k \right) \left( \sum_k p_k b_k \right) \geq \left( \sum_k p_k \sqrt{a_k b_k} \right)^2,$$

where  $a_k, b_k \geq 0$ .

# Uncertainty with the variance and the QFI

- Hence, we arrive at

$$\left[ \sum_k p_k (\Delta A)_{\varrho_k}^2 \right] \left[ \sum_k p_k (\Delta B)_{\varrho_k}^2 \right] \geq \frac{1}{4} \left[ \sum_k p_k L_{\varrho_k} \right]^2.$$

- We can choose the decomposition such that

$$\sum_k p_k (\Delta B)_{\varrho_k}^2 = F_Q[\varrho, B]/4.$$

- Due to the concavity of the variance we also know that

$$\sum_k p_k (\Delta A)_{\varrho_k}^2 \leq (\Delta A)^2.$$

- Hence, it follows that

$$(\Delta A)^2_{\varrho} F_Q[\varrho, B] \geq \left( \sum_k p_k L_{\varrho_k} \right)^2.$$

- In order to use the previous inequality, we need to know the decomposition  $\{p_k, \varrho_k\}$  that minimizes  $\sum_k p_k (\Delta B)_{\varrho_k}^2$ .

# Uncertainty with the variance and the QFI II

- We can have an inequality where we do not need to know that decomposition

$$(\Delta A)^2_{\varrho} F_Q[\varrho, B] \geq \left( \min_{\{p_k, \varrho_k\}} \sum_k p_k L_{\varrho_k} \right)^2.$$

- On the right-hand side, the bound is defined based on a convex roof.
- It can be shown that we can move to pure state decompositions.
- We know that

$$L_{\psi_k} \geq |\langle C \rangle_{\psi_k}|$$

holds.

# Uncertainty with the variance and the QFI II

- Then, we can obtain the inequality

$$(\Delta A)_{\varrho}^2 F_Q[\varrho, B] \geq \left( \min_{\{p_k, |\psi_k\rangle\}} \sum_k p_k |\langle C \rangle_{\psi_k}| \right)^2,$$

- Using

$$\sum_k p_k |\langle C \rangle_{\psi_k}| \geq \left| \sum_k p_k \langle C \rangle_{\psi_k} \right| \equiv |\langle C \rangle_{\varrho}|,$$

we arrive at the improved Heisenberg-Robertson uncertainty

$$(\Delta A)_{\varrho}^2 F_Q[\varrho, B] \geq |\langle C \rangle_{\varrho}|^2.$$

# Uncertainty with the variance and the QFI III

- The Heisenberg uncertainty

$$(\Delta A)_{\varrho}^2 (\Delta B)_{\varrho}^2 \geq \frac{1}{4} |\langle i[A, B] \rangle_{\varrho}|^2.$$

- The improved Heisenberg uncertainty

$$(\Delta A)_{\varrho}^2 F_Q[\varrho, B] \geq |\langle i[A, B] \rangle_{\varrho}|^2.$$

- It has been derived originally with a different method in

F. Fröwis, R. Schmied, and N. Gisin, Phys. Rev. A 92, 012102 (2015).

# Conditions for saturation

- Conditions for saturating the relation with the simple bound

$$(\Delta A)_{\varrho}^2 F_Q[\varrho, B] \geq \left( \min_{\{p_k, \varrho_k\}} \sum_k p_k L_{\varrho_k} \right)^2 \geq |\langle C \rangle_{\varrho}|^2.$$

- We have to have equality on the right-hand side.
- The optimal decomposition can be made with pure components  $\psi_k$ . Then, for all  $k, l$  we must have

$$\frac{1}{2} \langle \{A, B\} \rangle_{\psi_k} - \langle A \rangle_{\psi_k} \langle B \rangle_{\psi_k} = 0,$$

$$(\Delta A)_{\psi_k}^2 = (\Delta A)_{\psi_l}^2,$$

$$(\Delta B)_{\psi_k}^2 = (\Delta B)_{\psi_l}^2,$$

$$|\langle C \rangle_{\psi_k}| = |\langle C \rangle_{\varrho}|,$$

etc.

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# Cramér-Rao bound

- Error propagation formula

$$(\Delta\theta)_A^2 = \frac{(\Delta A)^2}{|\partial_\theta \langle A \rangle|^2} = \frac{(\Delta A)^2}{|\langle C \rangle|^2}.$$

- Then the precision of the estimation is bounded as

$$(\Delta\theta)^2 \geq \frac{1}{m} \min_A (\Delta\theta)_A^2,$$

where  $m$  is the number of independent repetitions.

- Cramér-Rao bound based on a convex roof

$$(\Delta\theta)^2 \geq \frac{1}{m} \times \underbrace{\frac{1}{4 \min_{\{p_k, |\psi_k\rangle\}} \left[ \sum_k p_k (\Delta B)^2_{\psi_k} \right]}}_{F_Q[\varrho, B]}.$$

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# Uncertainty relation based on a concave roof

- For any decomposition  $\{\rho_k, \varrho_k\}$  we have

$$(\Delta A)^2(\Delta B)^2 \geq \frac{1}{4} \left( \sum_k \rho_k L_{\varrho_k} \right)^2,$$

where

$$L_{\varrho} = \sqrt{|\langle \{A, B\} \rangle_{\varrho} - 2\langle A \rangle_{\varrho} \langle B \rangle_{\varrho}|^2 + |\langle C \rangle_{\varrho}|^2}.$$

- We can even take a concave roof on the right-hand side

$$(\Delta A)^2_{\varrho} (\Delta B)^2_{\varrho} \geq \frac{1}{4} \left( \max_{\{\rho_k, \varrho_k\}} \sum_k \rho_k L_{\varrho_k} \right)^2.$$

- We prove that for qubits the above inequality is saturated for all states.

# Any decomposition leads to a valid bound

- A simple inequality that is valid

$$(\Delta A)^2_{\varrho}(\Delta B)^2_{\varrho} \geq \frac{1}{4} \left( \sum_k \lambda_k L|k\rangle \right)^2,$$

if we have an eigendecomposition

$$\varrho = \sum_k \lambda_k |k\rangle \langle k|.$$

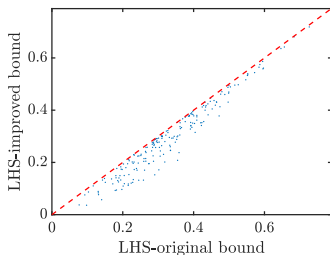
- We can even look for concave roof numerically.

# Numerical example

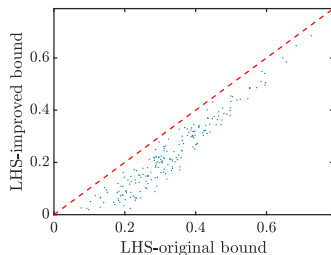
- For  $d = 3$

$$(\Delta J_x)_{\varrho}^2 (\Delta J_y)_{\varrho}^2 \geq \frac{1}{4} \left( \max_{\{p_k, \varrho_k\}} \sum_k p_k L_{\varrho_k} \right)^2.$$

- Eigenvalues  $J_x$  and  $J_y$  are  $-1, 0, +1$ .



Using the eigendecomposition



Numerical search

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# Uncertainty relations with a variance and the QFI

- Similar ideas work even for a sum of two variances. For example, for a continuous variable system

$$(\Delta x)^2 + (\Delta p)^2 \geq 1$$

holds, where  $x$  and  $p$  are the position and momentum operators.

- Hence, for any decompositions of the density matrix it follows that

$$\sum_k p_k (\Delta x)^2_{\psi_k} + \sum_k p_k (\Delta p)^2_{\psi_k} \geq 1.$$

- For  $p$  we choose the decomposition that leads to the minimal value for the average variance, i.e., the QFI over four.
- Then, since  $\sum_k p_k (\Delta x)^2_{\psi_k} \leq (\Delta x)^2$  holds, it follows that

$$(\Delta x)^2 + \frac{1}{4} F_Q[\varrho, p] \geq 1.$$

# Uncertainty relations with two variances and the QFI

- Let us start from the relation for pure states

$$(\Delta J_x)^2 + (\Delta J_y)^2 + (\Delta J_z)^2 \geq j,$$

where  $J_l$  are the spin components fulfilling

$$J_x^2 + J_y^2 + J_z^2 = j(j+1)\mathbb{1}.$$

- Based on similar ideas we arriving at

$$(\Delta J_x)^2 + (\Delta J_y)^2 + \frac{1}{4}F_Q[\varrho, J_z] \geq j.$$

- See parallel publication in



# Uncertainty relations with two variances and the QFI II

- For a spin- $j$  particle, the following inequality bounds from below the metrological usefulness of the state

$$F_Q[\varrho, J_z] \geq 4j - 4(\Delta J_x)^2 - 4(\Delta J_y)^2 =: B_{FQ}.$$

- For instance, this can be used in BEC of two-state atoms to bound  $F_Q[\varrho, J_z]$  with the variances.

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# Simple observation to prove further relations

- Let us consider a relation

$$\underbrace{(\Delta A)^2}_{\text{variance}} \rho \geq \underbrace{g(\rho)}_{\text{convex in } \rho} ,$$

which is true for pure states.

- If  $g(\rho)$  is convex in density matrices, then

$$\frac{1}{4} F_Q[\rho, A] \geq g(\rho)$$

holds for mixed states.

- Proof.*  $\frac{1}{4} F_Q[\rho, A]$  is given as a convex roof of the variance.
- It is the largest convex function that equals  $(\Delta A)^2_\rho$  for all pure states.



# Repeating the proof for the two variances and the QFI

- We rewrite the relation with three variances as

$$\underbrace{(\Delta J_x)^2}_{\text{variance}} \geq j - \underbrace{(\Delta J_y)^2 - (\Delta J_z)^2}_{\text{convex in } \varrho},$$

- The right-hand side is convex in  $\varrho$  and the left-hand side is a variance.
- Hence,

$$\frac{1}{4} F_Q[\varrho, J_z] \geq j - (\Delta J_x)^2 - (\Delta J_y)^2.$$

# Extreme spin squeezing

- For a particle with spin- $j$

$$\underbrace{(\Delta J_x)^2}_{\text{variance}} \geq \underbrace{j F_j(\langle J_z \rangle / j)}_{\text{convex in } \varrho}$$

holds, where  $F_j(X)$  is a convex function defined as

$$F_j(X) = \min_{\varrho: \langle J_z \rangle = Xj} \frac{(\Delta J_x)^2}{j}.$$

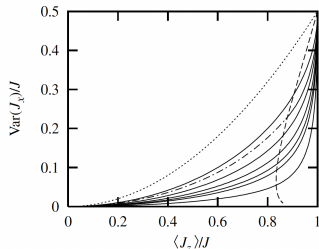


FIG. 1. Maximal squeezing for different values of  $J$ . The curves starting at the origin represent the minimum obtainable variance as a function of the mean spin. Starting from above, the curves represent  $J = 1/2, 1, 3/2, 2, 3, 4, 5,$  and  $10$ . The dotted curve for  $J = 1/2$  is the limit identified in Ref. [11]. The solid

# Application for extreme spin squeezing

- The metrological usefulness of a state is bounded with the spin-length as

$$F_Q[\varrho, J_x] \geq 4jF_j(\langle J_z \rangle / j).$$

- *Proof.* For the components of the angular momentum

$$(\Delta J_x)^2 \geq jF_j(\langle J_z \rangle / j)$$

holds. Then, it follows

$$\frac{1}{4} F_Q[\varrho, J_x] \geq jF_j(\langle J_z \rangle / j).$$



Can also be seen based on I. Apellaniz, M. Kleinmann, O. Gühne, and G. Tóth, Phys. Rev. A 95, 032330 (2017).

## Application for extreme spin squeezing II

- It is clear that if  $\langle J_z \rangle > 0$  then  $F_j(\langle J_z \rangle / j) > 0$ . Hence, it follows that  $F_Q[\varrho, J_x] > 0$ .
- Thus, if the z-component of the angular momentum has a non-zero expectation value then the state can be used for metrology with the Hamiltonian  $J_x$ .

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# CV systems

- Bosonic mode, canonical  $x$  and  $p$  operators. For coherent states,  $|\alpha\rangle$

$$(\Delta x)^2 = (\Delta p)^2 = \frac{1}{2}$$

holds. For mixtures of coherent states

$$\varrho_{\text{mc}} = \sum_k p_k |\alpha_k\rangle \langle \alpha_k|$$

due to the concavity of the variance and the convexity of the QFI

$$(\Delta x)^2, (\Delta p)^2 \geq \frac{1}{2}, \quad F_Q[x, \varrho], F_Q[p, \varrho] \leq 2.$$

# CV systems II

- For a mixture of products of coherent states  $\alpha_k^{(l)}$  of the form

$$\varrho_{\text{sepc}} = \sum_k p_k |\alpha_k^{(1)}\rangle\langle\alpha_k^{(1)}| \otimes |\alpha_k^{(2)}\rangle\langle\alpha_k^{(2)}|$$

for the the collective quantities

$$[\Delta(x_1 \pm x_2)]^2 \geq 1; \quad [\Delta(p_1 \pm p_2)]^2 \geq 1.$$

Moreover,

$$F_Q[\varrho, p_1 \pm p_2] \leq 4; \quad F_Q[\varrho, x_1 \pm x_2] \leq 4.$$

For such states the multi-variable Glauber-Sudarshan  $P$  function is non-negative.

- Consider entanglement detection in two-mode systems with uncertainty relations.
- A well-known entanglement criterion is

$$[\Delta(x_1 + x_2)]^2 + [\Delta(p_1 - p_2)]^2 \geq 2.$$

If a quantum state violates it, then it is entangled.

L.-M. Duan, G. Giedke, J. I. Cirac, and P. Zoller, *Phys. Rev. Lett.* 84, 2722 (2000);  
R. Simon, *Phys. Rev. Lett.* 84, 2726 (2000).

# CV systems IV

- For a two-mode state, the following uncertainty relation holds

$$\begin{aligned} [\Delta(x_1 + x_2)]^2 + [\Delta(p_1 - p_2)]^2 \geq \\ 4/F_Q[\varrho, p_1 + p_2] + 4/F_Q[\varrho, x_1 - x_2]. \end{aligned}$$

As a consequence, states violating the entanglement condition are metrologically more useful than states that are the mixtures of products of coherent states.

- *Proof.* We start from the relations

$$\begin{aligned} [\Delta(x_1 + x_2)]^2 F_Q[\varrho, p_1 + p_2] &\geq 4, \\ [\Delta(p_1 - p_2)]^2 F_Q[\varrho, x_1 - x_2] &\geq 4. \end{aligned}$$

- Then, in both inequalities we divide by the term containing the QFI. Finally, we sum the two resulting inequalities.

# CV systems V

- The violation of the entanglement criterion given implies the violation of one of the inequalities for the QFI.
- Thus, violation of the uncertainty relation-based entanglement condition also means that the state has larger metrological usefulness than mixtures of products of coherent states. ■
- We did not prove that violating the entanglement condition leads to larger metrological usefulness than that of separable states in general.
- Even for pure product states  $F_Q[\varrho, x_1 \pm x_2]$  or  $F_Q[\varrho, p_1 \pm p_2]$  can be arbitrarily large for two bosonic modes.

# Summary

- We showed how to derive new uncertainty relations with the variance and the quantum Fisher information based on simple convexity arguments.

See:

Géza Tóth and Florian Fröwis,

Uncertainty relations with the variance and the quantum Fisher information based on convex decompositions of density matrices,

[Phys. Rev. Research 4, 013075 \(2022\).](#)

THANK YOU FOR YOUR ATTENTION!