

Lower bounds on the quantum Fisher information based on the variance and various types of entropies

G. Tóth^{1,2,3}

¹Theoretical Physics, University of the Basque Country (UPV/EHU), Bilbao, Spain

²IKERBASQUE, Basque Foundation for Science, Bilbao, Spain

³Wigner Research Centre for Physics, Budapest, Hungary

DFG, Erlangen

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1 Motivation

- Estimating the quantum Fisher information is important

2 Background

- Quantum Fisher information
- Estimation of the QFI based on measurements

3 Results

- Bounding the quantum Fisher information based on the variance

Estimating the quantum Fisher information is important

- Many experiments are aiming to carry out a metrological task.
- The quantum Fisher information tells us the best precision achievable.

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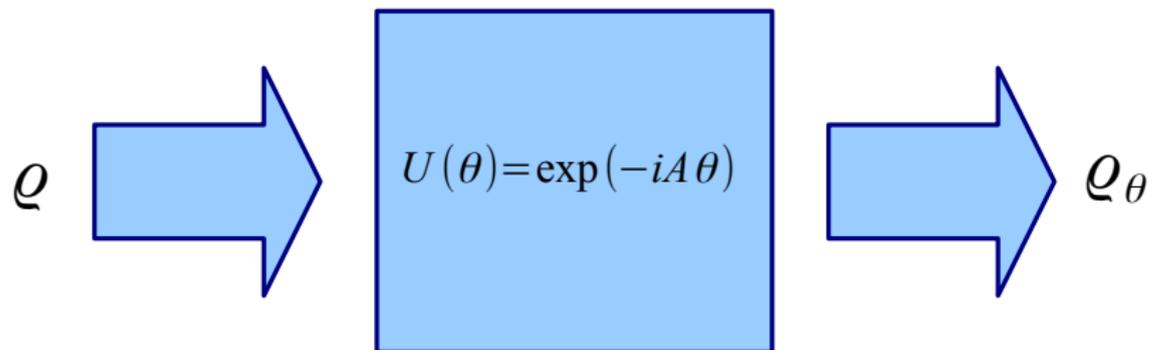
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Quantum metrology

- Fundamental task in metrology



- We have to estimate θ in the dynamics

$$U = \exp(-iA\theta).$$

The quantum Fisher information

- Cramér-Rao bound on the precision of parameter estimation

$$(\Delta\theta)^2 \geq \frac{1}{F_Q[\varrho, \mathbf{A}]}, \quad (\Delta\theta)^{-2} \leq F_Q[\varrho, \mathbf{A}].$$

where $F_Q[\varrho, \mathbf{A}]$ is the **quantum Fisher information**.

- Large $F_Q \rightarrow$ High precision.
- The quantum Fisher information is

$$F_Q[\varrho, \mathbf{A}] = 2 \sum_{k,l} \frac{(\lambda_k - \lambda_l)^2}{\lambda_k + \lambda_l} |\langle k | \mathbf{A} | l \rangle|^2,$$

where $\varrho = \sum_k \lambda_k |k\rangle\langle k|$.

Properties of the Fisher information

Properties of the QFI

- For pure states, it equals four times the variance,
 $F_Q[|\Psi\rangle\langle\Psi|, A] = 4(\Delta A)^2_\Psi$.
- For mixed states, $F_Q[\varrho, A] \leq 4(\Delta A)^2_\varrho$.
- Convex in ϱ .

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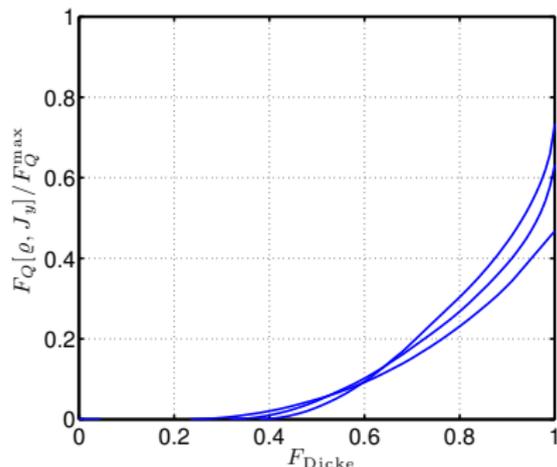
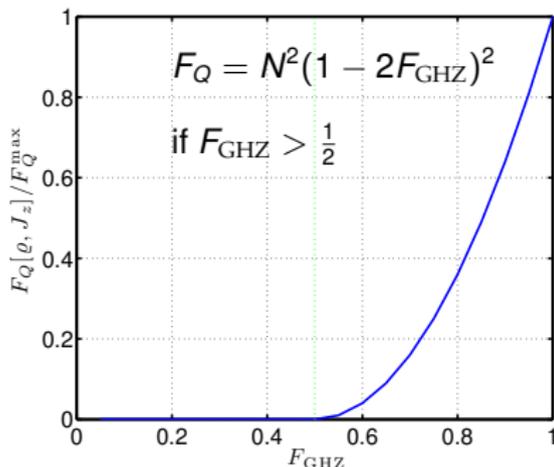
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Example: estimate QFI based on measurements

- **Lower** bound on the quantum Fisher information based on some measurements



Quantum Fisher information vs. Fidelity with respect to
(a) GHZ states and (b) Dicke states
for $N = 4, 6, 12$. $F_Q^{\max} = N^2$.

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Bound based on the variance

We define the quantity

$$V(\varrho, A) := (\Delta A)^2 - \frac{1}{4}F_Q[\varrho, A].$$

- Concave in ϱ , a generalized variance.
[Petz, J. Phys. A 35, 929 (2002);
Gibilisco, Hiai, and Petz, IEEE Trans. Inf. Theory 55, 439 (2009)]
- $V(\varrho, A) = 0$ for pure states.
- For states close to be pure, $V(\varrho, A)$ is small.
- For states that are far from pure, $V(\varrho, A)$ can be larger.

Bound based on the variance, rank-2

Observation 1.—For rank-2 states ϱ , the equality

$$(\Delta A)^2 - \frac{1}{4}F_Q[\varrho, A] = \frac{1}{2}[1 - \text{Tr}(\varrho^2)](\tilde{\sigma}_1 - \tilde{\sigma}_2)^2$$

holds, where $\tilde{\sigma}_k$ are the nonzero eigenvalues of the matrix

$$A_{kl} = \langle k|A|l\rangle.$$

Here $|k\rangle$ are the two eigenvectors of ϱ with nonzero eigenvalues.

Note

$$S_{\text{lin}}(\varrho) = 1 - \text{Tr}(\varrho^2) = 1 - \sum_k \lambda_k^2 = \sum_{k \neq l} \lambda_k \lambda_l.$$

Bound based on the variance, rank-2 II

- Two-dimensional subspace

$$\{|000..00\rangle, |111..11\rangle\}.$$

- Relevant for experiments with trapped ions, creating GHZ states

$$|\text{GHZ}\rangle = \frac{1}{\sqrt{2}} (|000..00\rangle + |111..11\rangle).$$

- The QFI is obtained as

$$F_Q[\varrho, J] = 2N^2 \left[\text{Tr}(\varrho^2) - \langle P_{000..00} \rangle^2 - \langle P_{111..11} \rangle^2 \right].$$

[GHZ experiments with ione traps: D. Leibfried *et al.*, Science 2004;
C. Sackett *et al.*, Nature 2000; T. Monz *et al.*, Phys. Rev. Lett. 2011.]

Bound based on the variance, arbitrary rank

Observation 2.—For states ϱ with an arbitrary rank we have

$$(\Delta A)^2 - \frac{1}{4}F_Q[\varrho, A] \leq \frac{1}{2}S_{\text{lin}}(\varrho)[\sigma_{\max}(A) - \sigma_{\min}(A)]^2,$$

where $\sigma_{\max}(A)$ and $\sigma_{\min}(A)$ are the largest and smallest eigenvalue of A , respectively.

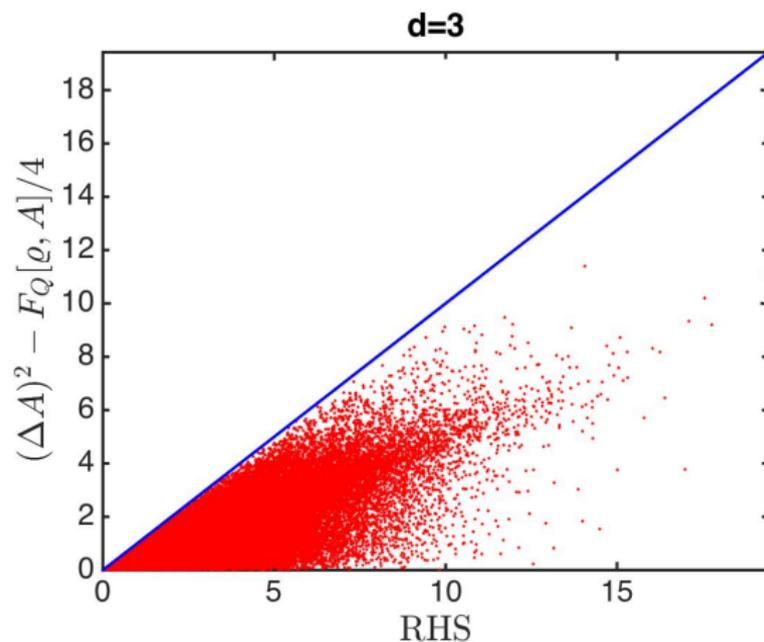
Estimate QFI:

$$F_Q[\varrho, A] \geq 4(\Delta A)^2 - 2S_{\text{lin}}(\varrho)[\sigma_{\max}(A) - \sigma_{\min}(A)]^2.$$

- 1 Measure the variance.
- 2 Estimate the purity.
- 3 Find a lower bound on F_Q .

Bound based on the variance, arbitrary rank II

Numerical verification of the bound



Quantities averaged over SU(d) generators

- Traceless Hermitian matrices

$$A_{\vec{n}} := \vec{A}^T \vec{n},$$

where $\vec{A} = [A^{(1)}, A^{(2)}, A^{(3)}, \dots]^T$ are the SU(d) generators.

- $N_g = d^2 - 1$ is the number of generators.
- Average over unit vectors

$$\text{avg}_{\vec{n}} (\Delta A_{\vec{n}})^2 = \int (\Delta A_{\vec{n}})^2 M(d\vec{n}) = \frac{1}{N_g} \sum_{k=1}^{N_g} (\Delta A_k)^2.$$

Similar statement holds for $F_Q[\varrho, A]$ and $V(\varrho, A)$.

Bound on the average V

Observation 3.—Average of V over traceless Hermitian matrices

$$\text{avg}_{\vec{n}} V(\varrho, A_{\vec{n}}) = \frac{2}{d^2 - 1} \left[S_{\text{lin}}(\varrho) + H(\varrho) - 1 \right],$$

d is the dimension of the system.

The quantity $H(\varrho)$ is defined as

$$H(\varrho) = 2 \sum_{k,l} \frac{\lambda_k \lambda_l}{\lambda_k + \lambda_l} = 1 + 2 \sum_{k \neq l} \frac{\lambda_k \lambda_l}{\lambda_k + \lambda_l}.$$

Average V is zero only for pure states. \rightarrow Similar to entropies.

Average quantum Fisher information

- Average of the quantum Fisher information

$$\text{avg}_{\vec{n}} F_Q[\varrho, \mathbf{A}_{\vec{n}}] = \frac{8}{N_g} [d - H(\rho)].$$

- Maximal for pure states.

Bound based on the QFI based on H

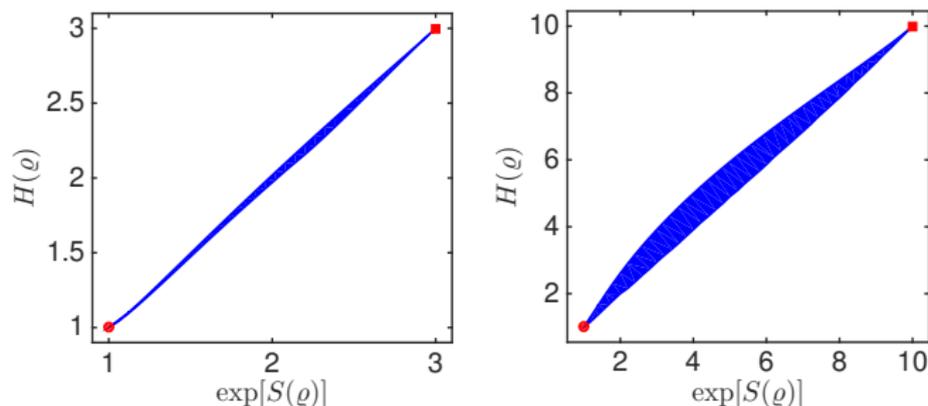


Figure: Relation of $H(\rho)$ and the von Neumann entropy for $d = 3$ and 10 .

- (filled area) Physical quantum states.
- (circle) Pure states.
- (square) Completely mixed state.

We see that

$$H(\rho) \sim \exp[S(\rho)]. \quad \text{Remember: } \text{avg } F_Q = \frac{8}{N_g} [d - H(\rho)].$$

Bound based on the QFI based on H

Message:

Large entropy \rightarrow Small average QFI

Small entropy \rightarrow Large average QFI

The more mixed the state,
the less useful it is for metrology.

Kubo-Mori-Bogoliubov quantum Fisher information I

Consider

$$F_Q^{\log}[\varrho, \mathbf{A}] = \sum_{k,l} [\log(\lambda_k) - \log(\lambda_l)] (\lambda_k - \lambda_l) |\mathbf{A}_{kl}|^2,$$

that fulfils

$$\frac{d^2}{d^2\theta} S(\varrho || e^{-i\mathbf{A}\theta} \varrho e^{i\mathbf{A}\theta})|_{\theta=0} = F_Q^{\log}[\varrho, \mathbf{A}].$$

We found

$$\text{avg}_{\vec{n}} F_Q^{\log}[\varrho, \mathbf{A}_{\vec{n}}] = -\frac{2}{N_g} (2dS + 2 \sum_k \log \lambda_k).$$

Kubo-Mori-Bogoliubov quantum Fisher information II

- Relation to other works in the literature:

Quantum version of the classical isoperimetric inequality relating the KMB QFI and the $\exp(S)$ for Gaussian states.

[S. Huber, R. Koenig, and A. Vershynina, arxiv:1606.08603;
C. Rouze, N. Datta, and Y. Pautrat, arxiv:1607.04242.]

Summary

- We discussed how to find lower bounds on the quantum Fisher information with the variance and the entropy.

See:
G. Tóth,

Lower bounds on the quantum Fisher information based on the variance and various types of entropies,
[arxiv:1701.07461](https://arxiv.org/abs/1701.07461).

THANK YOU FOR YOUR ATTENTION!

